

B-physics theory

Thomas Becher, Fermilab

The Tevatron Connection,
June 24, 2005

The Tevatron Connection

A Series of CDF & DØ Presentations with Theoretical Perspectives

JUNE 24 - 25, 2005

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FERMILAB, I

$$\begin{aligned} \mathcal{L}_{EW} = & \bar{R} i \gamma^\mu (\partial_\mu + i g' A_\mu Y) R + \bar{L} i \gamma^\mu (\partial_\mu + i g' A_\mu Y + i g \vec{E} \cdot \vec{b}_\mu) L \\ & - \frac{1}{4} \vec{F}_{\mu\nu} \cdot \vec{F}^{\mu\nu} - \frac{1}{4} f_{\mu\nu} f^{\mu\nu} + (\partial^\mu \varphi)^\dagger (\partial_\mu \varphi) - V(\varphi^\dagger \varphi) \\ & - \sum_{ij} S_{ij} [(\bar{L}_i \bar{\varphi}) q_{iR} + \bar{q}_{iR} (\varphi^\dagger L_j)] \end{aligned}$$

$$\begin{aligned} \mathcal{L}_{EW} = & \bar{R} i \gamma^\mu (\partial_\mu + i g' A_\mu Y) \\ & - \frac{1}{4} \vec{F}_{\mu\nu} \cdot \vec{F}^{\mu\nu} \\ & - \sum_{ij} S_{ij} \end{aligned}$$

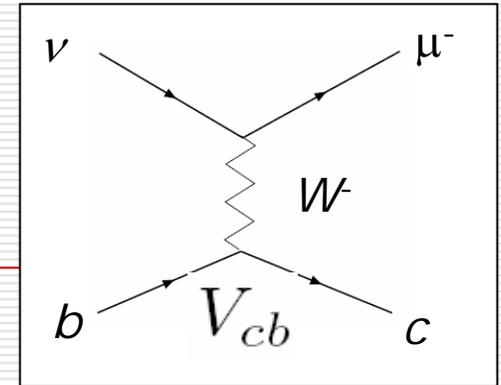
Yukawa couplings

- Fermion masses
 - Flavor changing interactions
 - CP violation
- } CKM matrix

Overview

- CKM status and news
 - Exclusive V_{ub} from the lattice
 - $\sin(2\beta)$: trees vs. penguins
- B-physics with “collider physics methods”
 - Heavy-to-light decays: Factorization, Soft-Collinear Effective Theory, ...
 - Scope, limitations and relevance for B-physics program at the Tevatron

The CKM matrix



$$V_{\text{CKM}} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

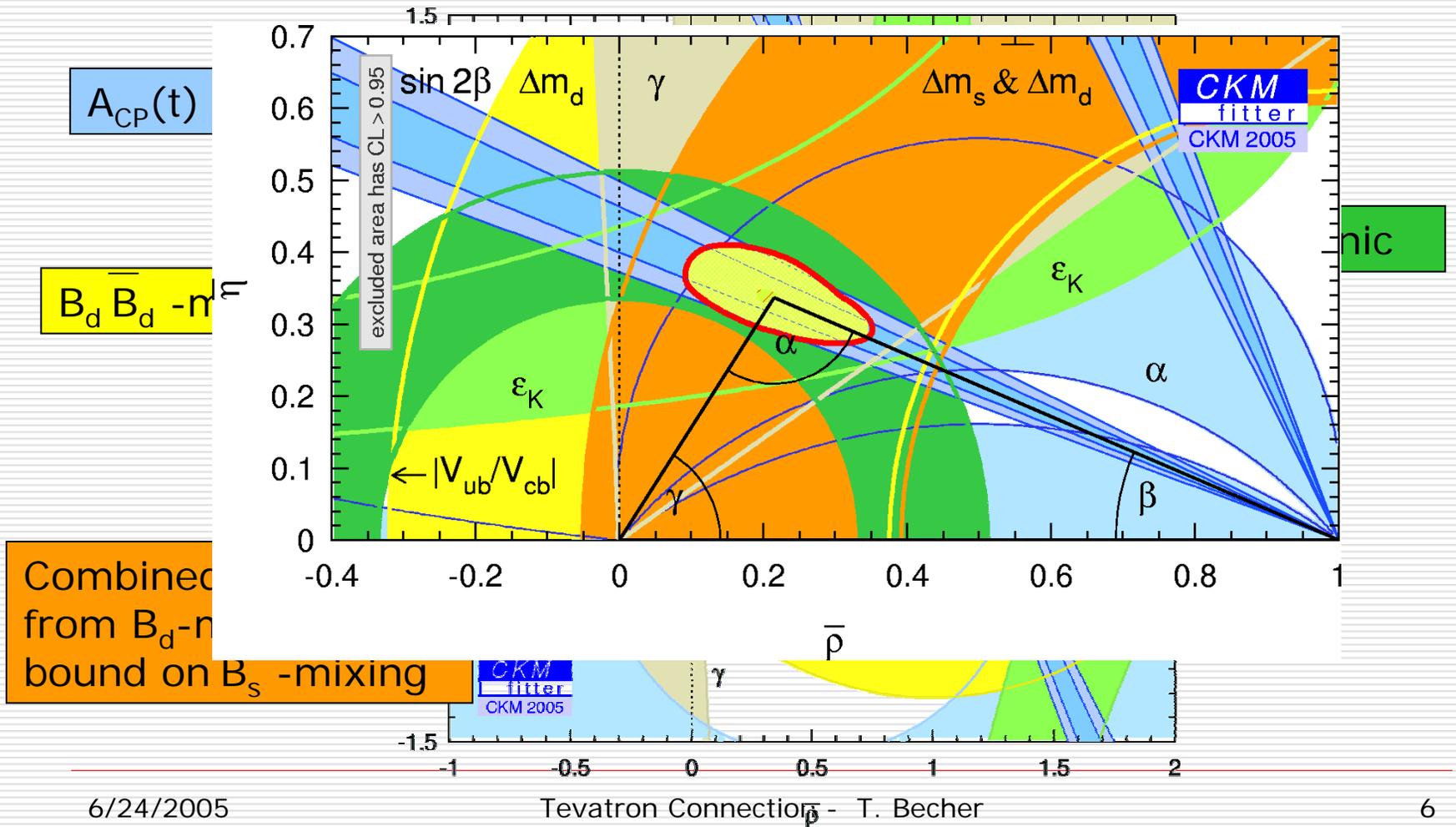
- Flavor changing quark interactions
- Unitary. 4 parameters including 1 complex CP violating phase

Wolfenstein parameterization

$$V_{\text{CKM}} = \begin{pmatrix} 1 - \frac{\lambda^2}{2} & \lambda & A \lambda^3 (\rho - i \eta) \\ -\lambda & 1 - \frac{\lambda^2}{2} & A \lambda^2 \\ A \lambda^3 (1 - \rho - i \eta) & -A \lambda^2 & 1 \end{pmatrix}$$

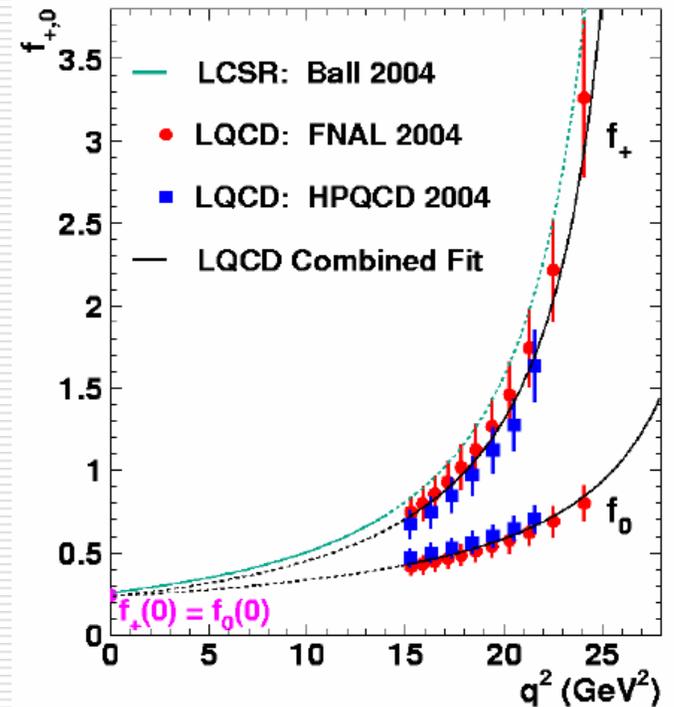
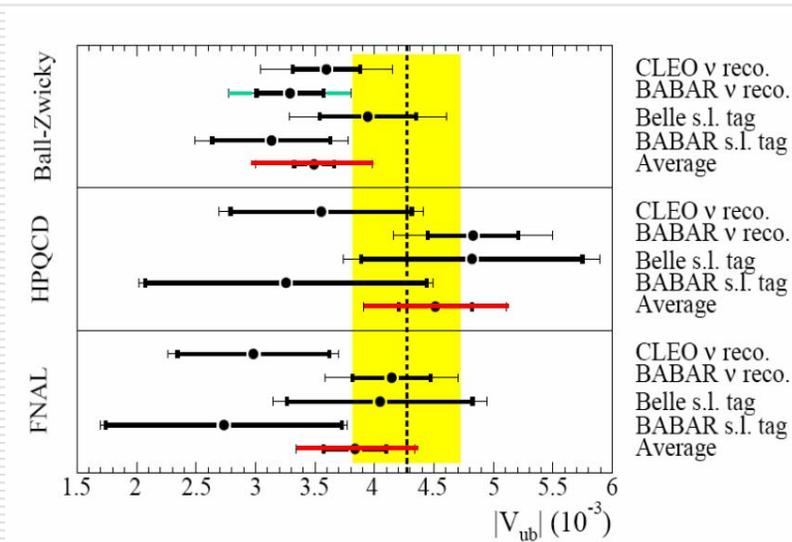
- $\lambda = |V_{us}| \approx 0.22$. ($A \approx 0.8$, $\rho \approx 0.2$, $\eta \approx 0.3$)
- Almost diagonal. Hierarchical.

Constraints on ρ, η : the unitarity triangle

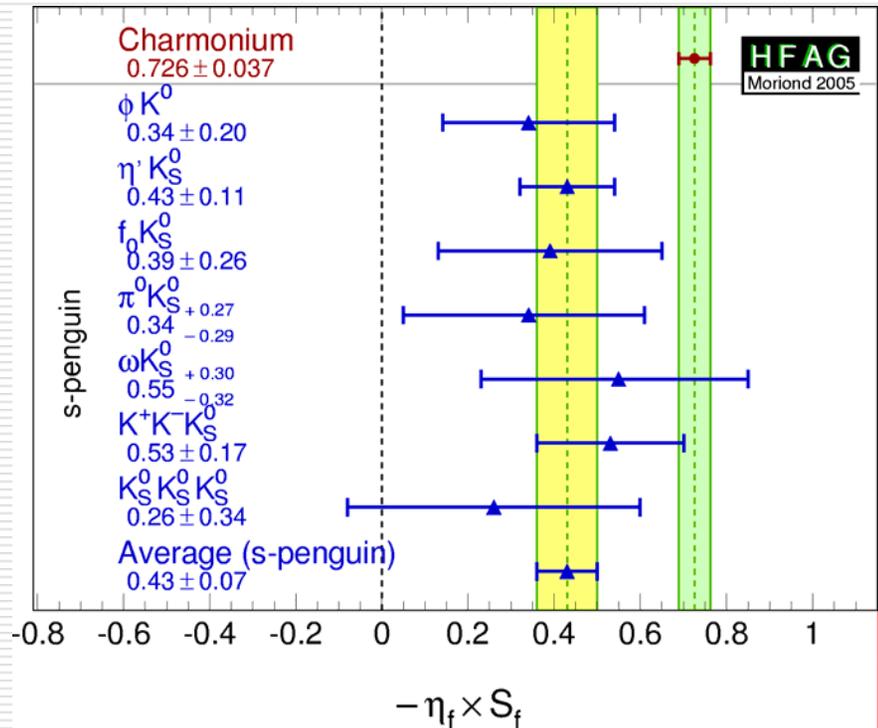
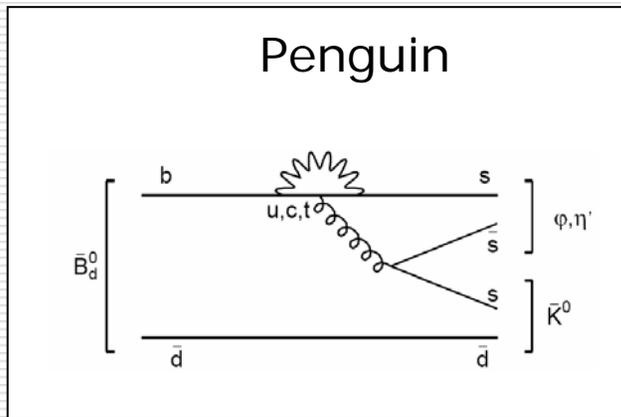
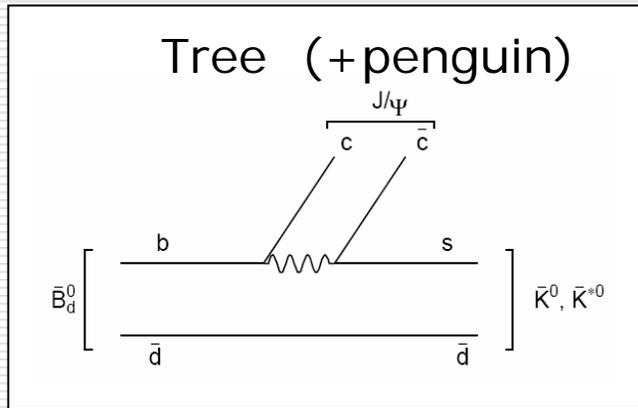


Exclusive $|V_{ub}|$ from $B \rightarrow \pi/\rho$

- **New:** *Dynamical* ($n_f=3$) lattice calculations of semileptonic form factors.



sin(2β) from loop dominated processes



Questions

□ Babar vs. Belle?

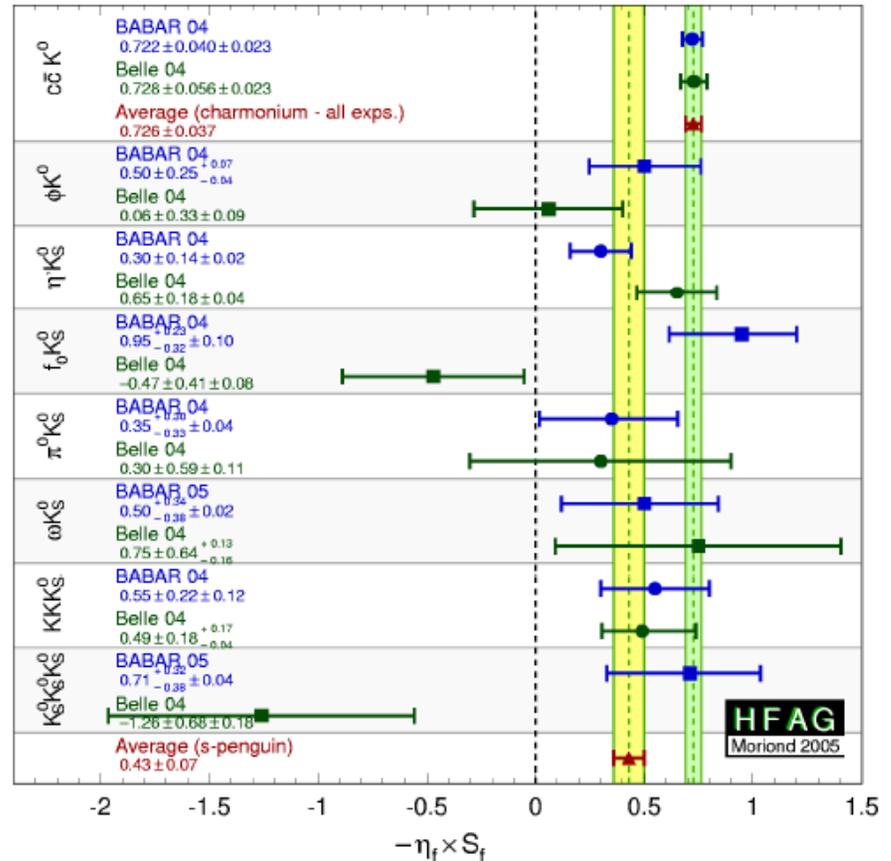
□ Theoretical uncertainties?

Contribution from small amplitudes?

■ Suppressed by

$$\left| \frac{V_{ub} V_{us}^*}{V_{cb} V_{cs}^*} \right| \approx \frac{1}{50}$$

■ Calculable in QCDF/SCET



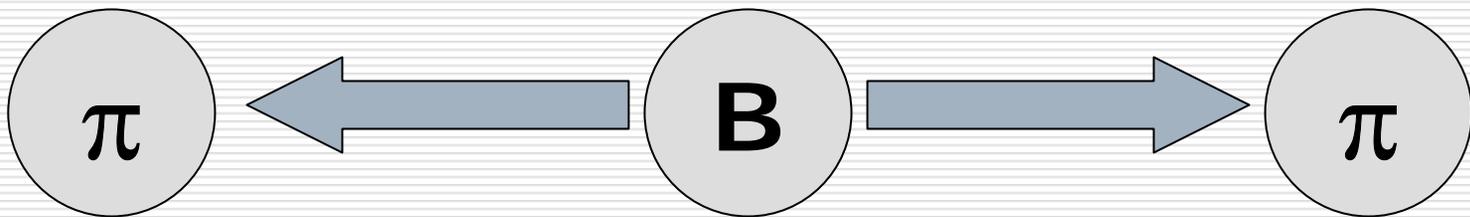
Collider physics methods in B-physics?

Factorization, Soft and Collinear particles, Jet-functions, Distribution Amplitudes...

WHAT IS WRONG WITH THE B-THEORISTS?

Hard scattering for $m_b \rightarrow \infty$

- In heavy-to-light decays, the decay products are typically energetic, e.g.



$$E_\pi = \frac{m_B}{2}$$

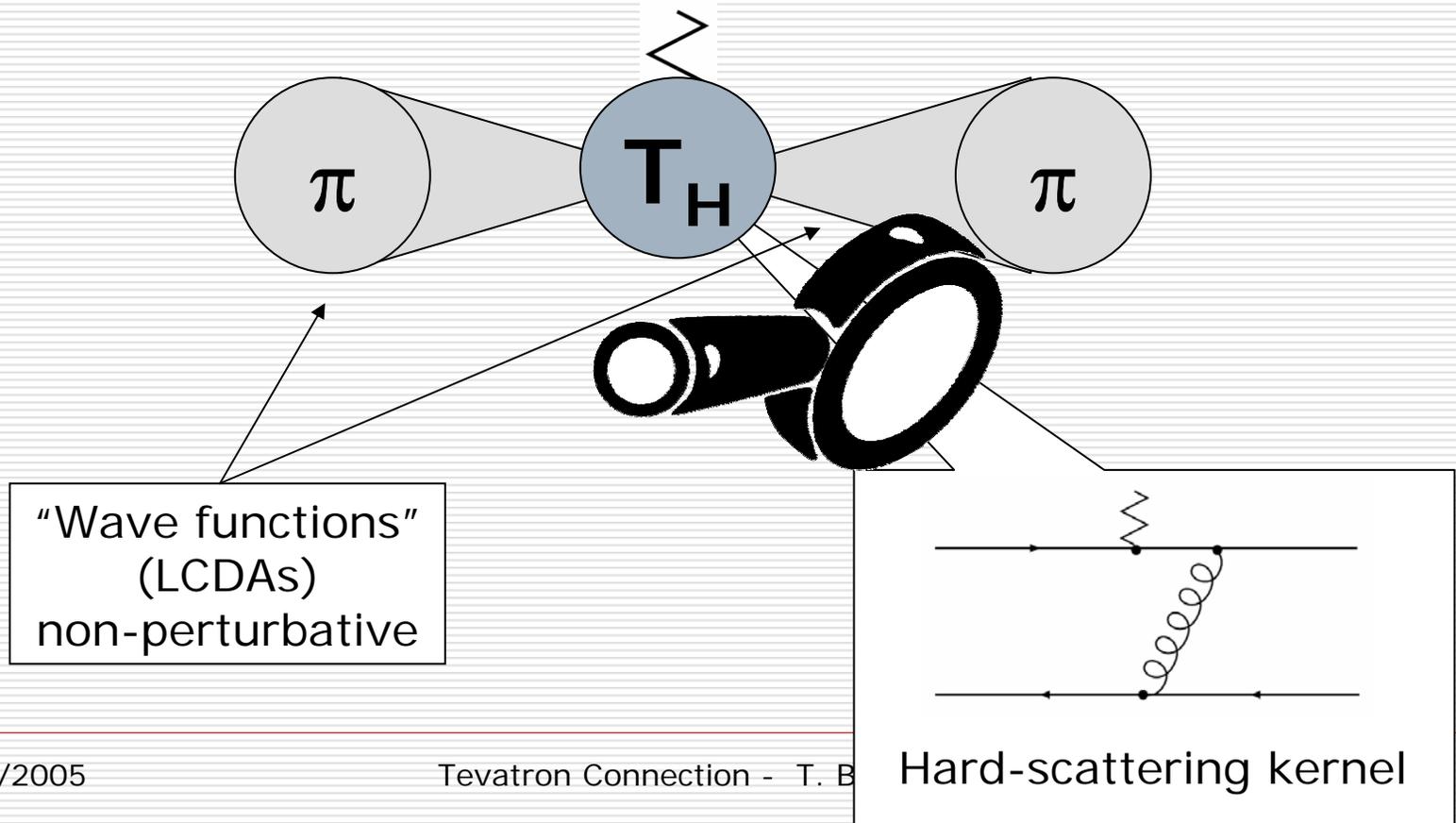
$$E_\pi = \frac{m_B}{2}$$

- Large momentum transfer \rightarrow factorization?

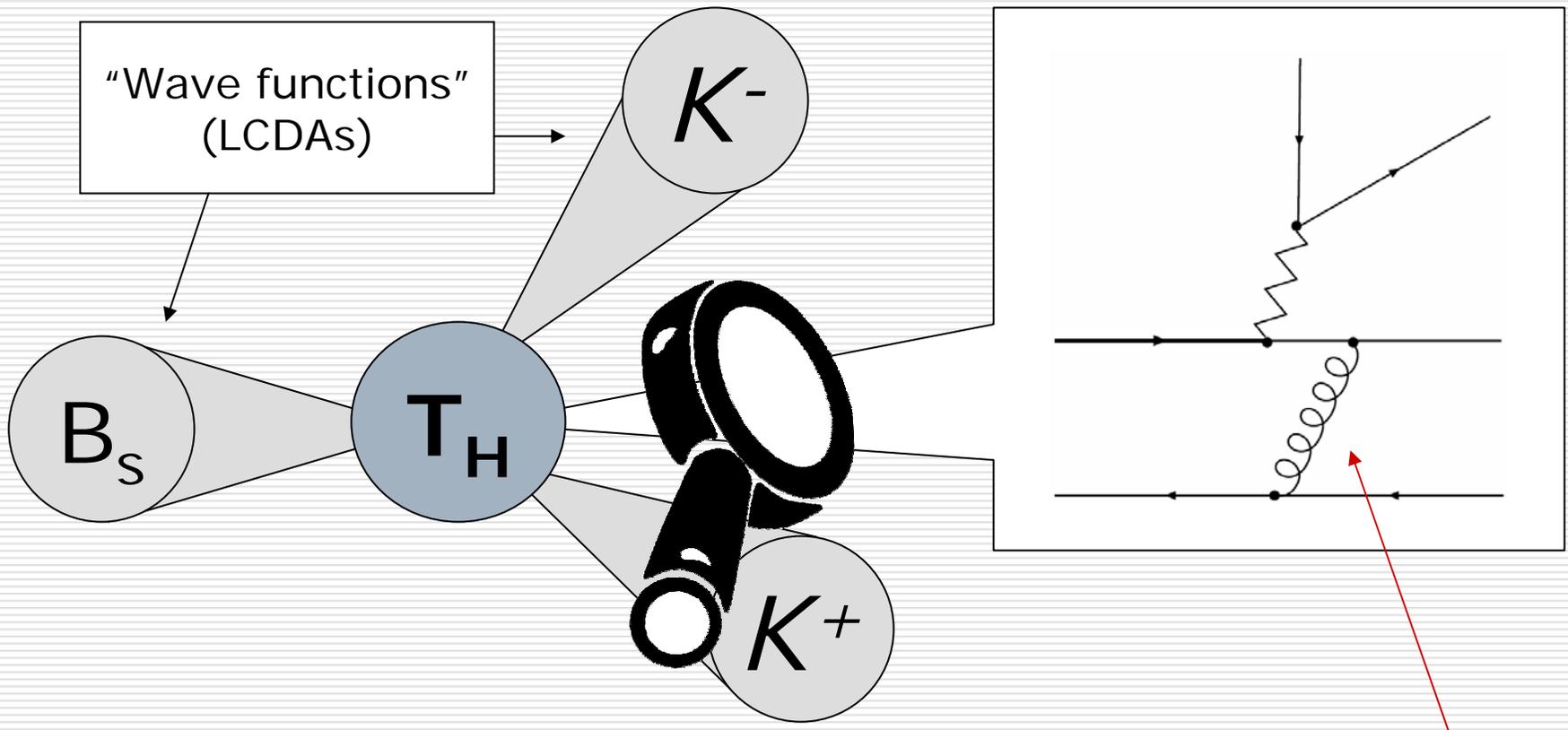
Remember?

(1980)

□ $\pi\pi$ form factor at large Q^2



Factorization for $m_b \rightarrow \infty$?

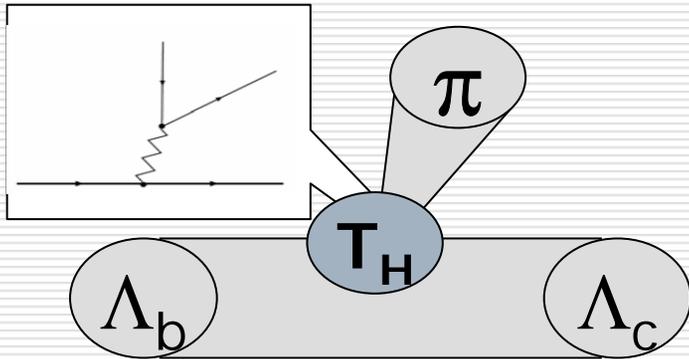


Factorization FAILS!

Ways out?

- Ignore the problem: Naïve factorization.
- Assume that non-factorizable effects are small after Sudakov resummation: “pQCD”.
- Give up on complete factorization. Allow for non-factorizable form factor piece: “QCD factorization”.

Example: $\Lambda_b \rightarrow \Lambda_c \pi$



Form factor (from semileptonic decay)

$$= 1 + \mathcal{O}(\alpha_s)$$

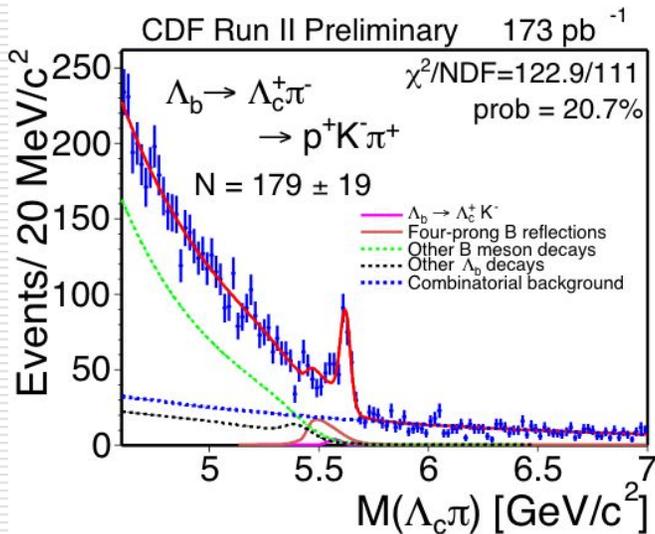
$$\mathcal{A} \propto V_{cb} V_{ud}^* \mathcal{F}_{\Lambda_b \rightarrow \Lambda_c}(q^2 \approx 0) f_\pi \int_0^1 dx T_H(x) \phi_\pi(x) + \mathcal{O}(\Lambda/E_\pi)$$

- Note: Obtain naïve factorization result to lowest order in α_s .
- Decays to two light hadrons also contain a completely factorizable piece.

CDF:
$$\frac{\mathcal{B}(\Lambda_b \rightarrow \Lambda_c^+ \mu^- \bar{\nu}_\mu)}{\mathcal{B}(\Lambda_b \rightarrow \Lambda_c^+ \pi^-)} = 20.0 \pm 3.0 (stat) \pm 1.2 (syst) \begin{matrix} +0.7 \\ -2.1 \end{matrix} (BR) \pm 0.5 (UBR)$$

$$\mathcal{B}(\Lambda_b \rightarrow \Lambda_c^+ \pi^-) = \left(0.41 \pm 0.19 (stat \oplus syst) \begin{matrix} +0.06 \\ -0.08 \end{matrix} (P_T \text{ spectrum}) \right) \%$$

$$\mathcal{B}(\Lambda_b \rightarrow \Lambda_c^+ \mu^- \bar{\nu}_\mu) = \left(8.1 \pm 1.2 (stat) \begin{matrix} +1.1 \\ -1.6 \end{matrix} (syst) \pm 4.3 (\mathcal{B}(\Lambda_b \rightarrow \Lambda_c^+ \pi^-)) \right) \%$$



□ Theory prediction?

Leibovich et al. hep-ph/0312319

- Needs $F(q^2 \approx 0)$ as an input!

$$\mathcal{B}(\Lambda_b \rightarrow \Lambda_c^+ \pi^-) \approx 0.46\%$$

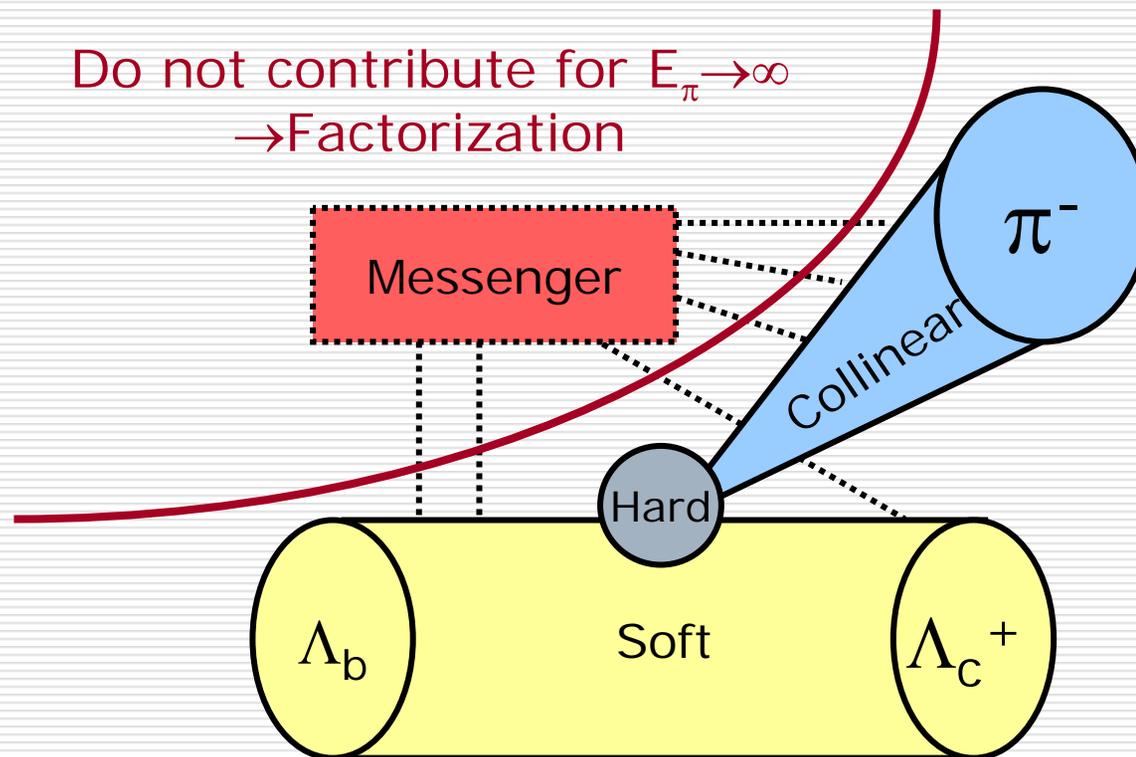
same $F(q^2)$ gives

$$\mathcal{B}(\Lambda_b \rightarrow \Lambda_c^+ \mu^- \bar{\nu}) \approx 6\%$$

Soft-Collinear Effective Theory

- Effective Lagrangian instead of diagrammatic language: systematic and transparent analysis of factorization properties.
- The “HQET” for heavy-to-light processes.
- Complicated field content
 - Soft-fields: partons of b- (and c-) hadrons
 - Collinear fields: partons of energetic light particles
 - Messenger fields: low energy interactions between soft and collinear partons. Absent for factorizable quantities.
- Hard scattering kernels $T_H \leftrightarrow$ Wilson coefficients of operators in SCET
- Based on “Strategy of Regions”. Fields in SCET \leftrightarrow Relevant momentum regions for expansion of diagrams.

$\Lambda_b \rightarrow \Lambda_c \pi$ analysis in SCET



Applications

- Factorization analysis (M : light meson)
 - $B_{s,d} \rightarrow D_{s,d} M, \Lambda_b \rightarrow \Lambda_c M$
 - $B_{s,d} \rightarrow M_1 M_2, B \rightarrow M l\nu$
 - $B_{s,d} \rightarrow M \gamma, B_{s,d} \rightarrow M \mu^+ \mu^-, B \rightarrow \gamma l\nu$
 - $B \rightarrow X_u l\nu, B \rightarrow X_s \gamma$ (end-point region)
- Sudakov resummation
 - $B \rightarrow M l\nu, B \rightarrow K^* \gamma, B \rightarrow \gamma l\nu$
 - $B \rightarrow X_u l\nu, B \rightarrow X_s \gamma$
- Λ/m_b power corrections
 - $B \rightarrow X_u l\nu, B \rightarrow X_s \gamma$ (“subleading shape functions”)
 - $B^0 \rightarrow D^{(*)0} \pi^0$
 - partial results for $B_{s,d} \rightarrow M_1 M_2$

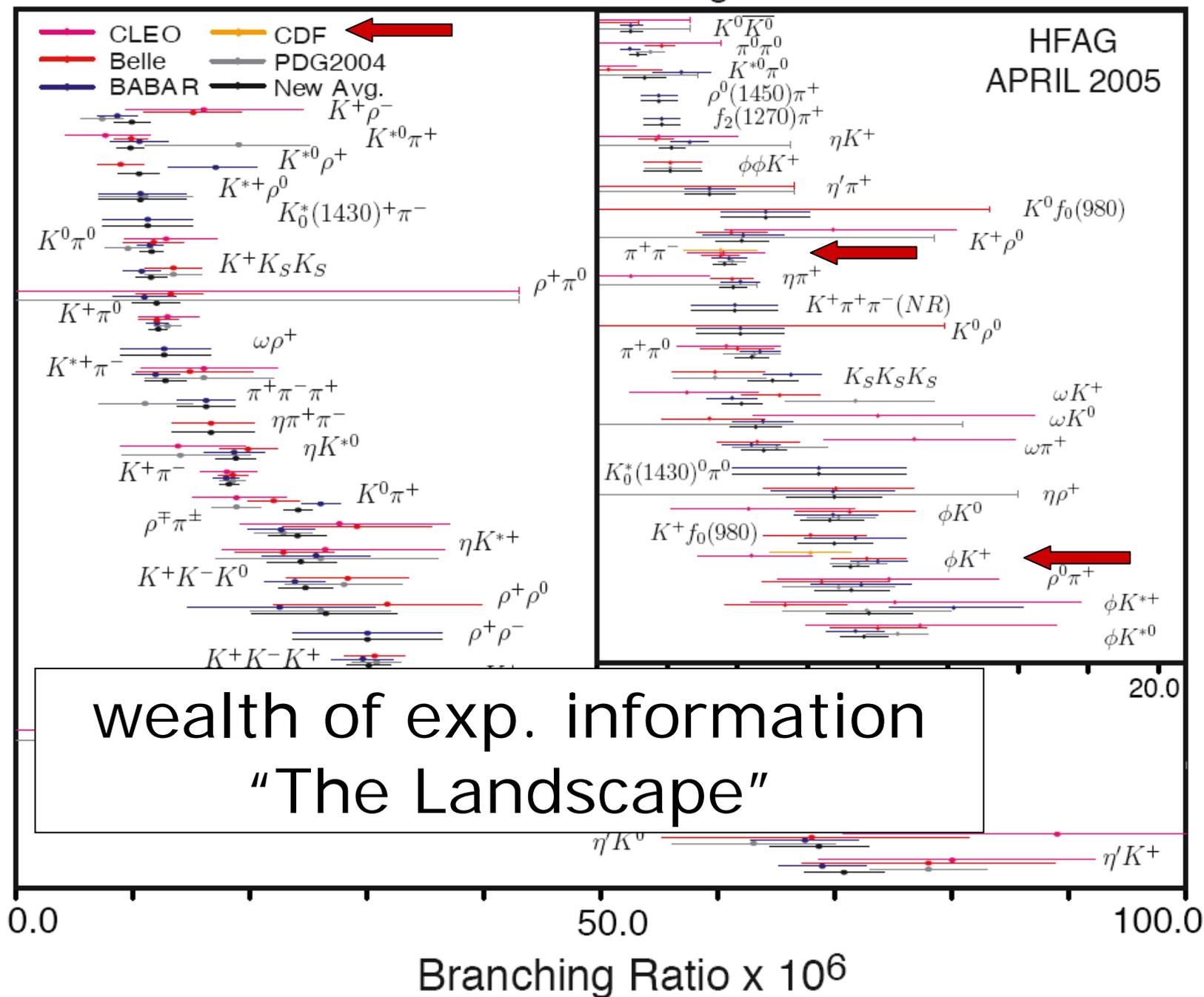
Hadronic input for two-body decays

- Need form-factors at $q^2=0$ as well as meson wave functions as input.
 - Have $B \rightarrow D$ form factor
 - Some knowledge on $B \rightarrow \pi$ form factors (exp., lattice)
 - Little lattice info on wave functions. B-meson wave function is especially difficult.
 - Can use SU(2) and SU(3) to relate hadronic parameters.
- Many factorization predictions rely on (light-cone) sum rule determinations for form factors and wave functions.

Strategies for charmless decays

- Beneke et al. "QCDF" ($B \rightarrow PP, B \rightarrow VP$)
 - Use sum rule input for leading power
 - Estimate enhanced power corrections
 - chiral enhancement, annihilation
- Bauer et al. "SCET" ($B \rightarrow \pi\pi$)
 - Use $B \rightarrow \pi\pi$ data to determine hadronic parameters.
 - Do not attempt to factorize "charming penguins"
 - To retain predictability: drop power and $\alpha_s(m_b)$ corrections.
- CKM fitter group: fit to charmless decays in QCDF.

Charmless B Branching Fractions



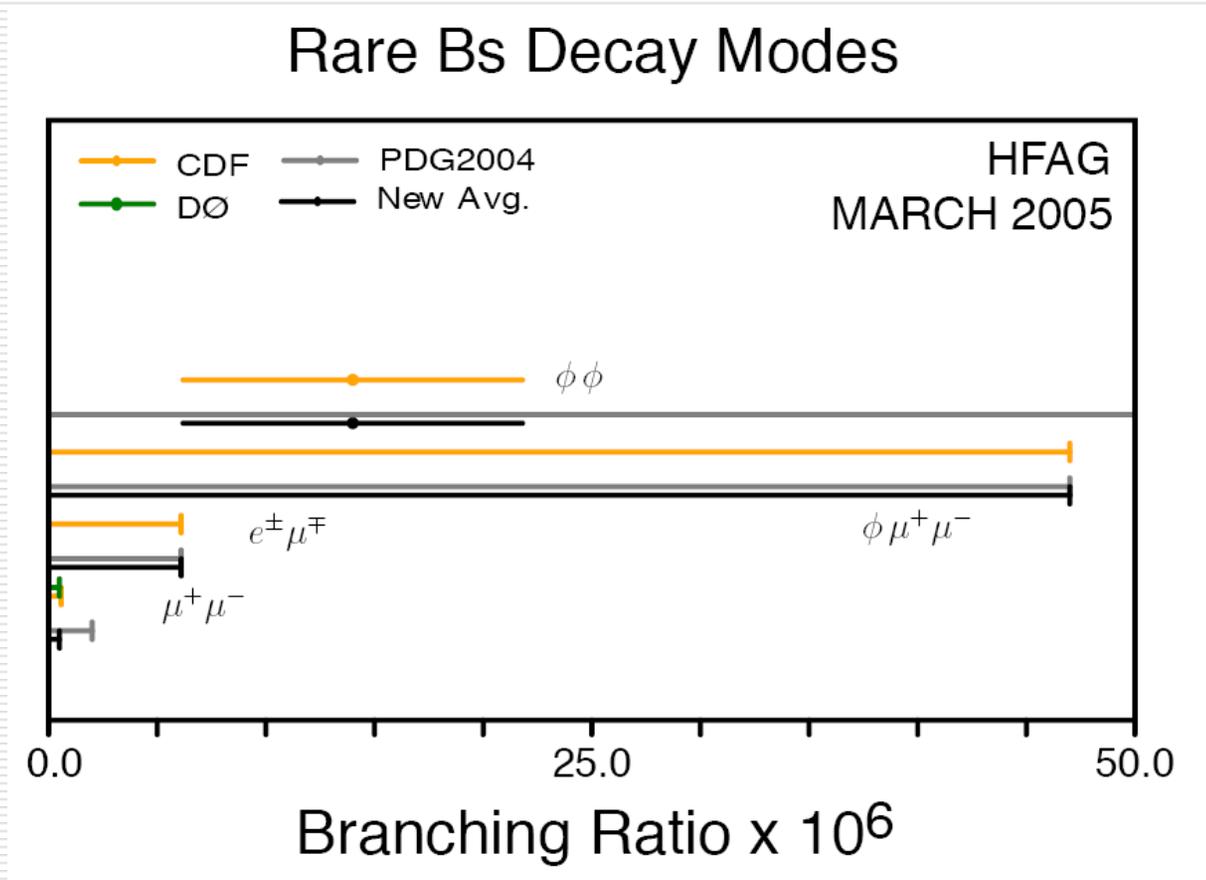
CDF results for $B_{s,d} \rightarrow hh$

- CDF sample of $B_d \rightarrow hh$ is as large as BaBar sample (and growing faster).
- Only a few decay channels measured up to now, but with *competitive accuracy*

Mode	BaBar	Belle	CDF
$10^6 \text{ Br}(B \rightarrow \pi^+ \pi^-)$	$4.7 \pm 0.6 \pm 0.2$	$4.4 \pm 0.6 \pm 0.3$	$4.5^{+1.4+0.5}_{-1.2-0.4}$
$10^6 \text{ Br}(B \rightarrow K^+ K^-)$	< 0.6	< 0.7	< 3.1
$10^6 \text{ Br}(B \rightarrow \phi K^+)$	$10^{+0.9}_{-0.8} \pm 0.5$	$9.60 \pm 0.92^{+1.05}_{-0.84}$	$7.6 \pm +1.3 \pm 0.6$
$A_{CP}(B \rightarrow \phi K^+)$	$0.054 \pm 0.056 \pm 0.012$	$0.01 \pm 0.12 \pm 0.05$	$-0.07 \pm 0.17^{+0.03}_{-0.02}$

- In future: $B_s \rightarrow KK, K\pi, \pi\pi, \dots$

The first B_s charmless two-body decay!



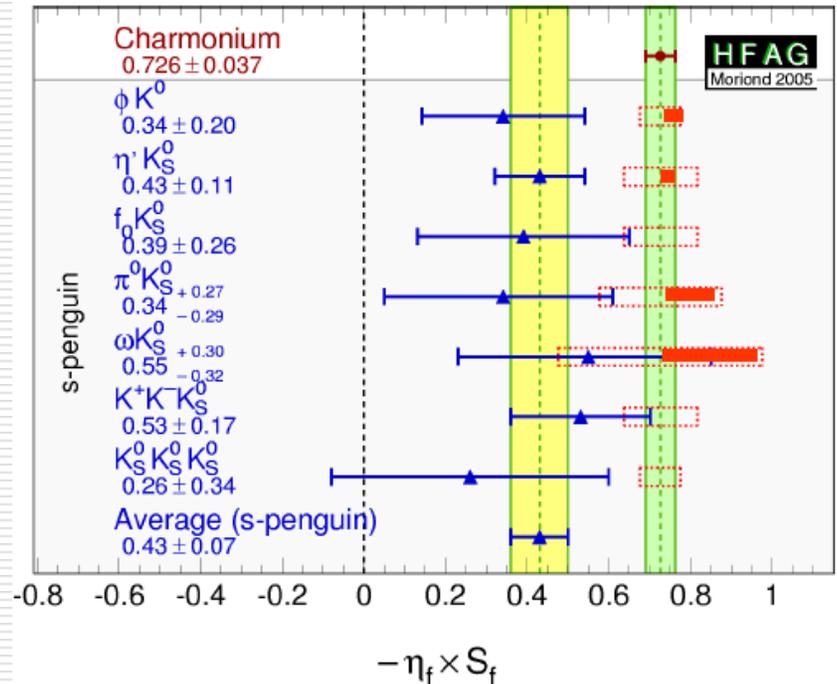
Outlook

(see talks of U. Nierste and R. Harr at INT workshop)

- $B_{d,s} \rightarrow h^+ h^-, B_s \rightarrow \phi\phi, \dots$
 - Updated $\text{Br}()$ results
 - more theoretical work on VV decays
 - $B_s \rightarrow \phi\phi, B_s \rightarrow K^+ K^-$ lifetime measurements?
 - Sensitive to *CP phase in $b \rightarrow \bar{q}qs$*
 - Isospin violating decay $B_s \rightarrow \rho^0 \phi$?
 - Time-dep. CP asymmetries?
- Baryonic decays $\Lambda_b \rightarrow p(\pi, K)$
 - Direct CP asymmetry
 - Factorization analysis (hadronic parameters?)
- Rare decay $B_s \rightarrow \phi \mu^+ \mu^-$
- γ from $B \rightarrow D^0 K$

sin(2β) from penguin dominated modes

- Beneke (hep-ph/0505075) evaluates the corrections to sin(2β) from the small amplitudes.
- Scans over all input parameters, (require BR's to 3σ). Gives resulting range.
- Positive contribution to ΔS. Strengthens the discrepancy!



Summary

- With the exception of $\sin(2\beta)$ from penguin dominated modes all measurements are consistent with the CKM picture of the flavor sector.
 - theoretical uncertainty on $\sin(2\beta)$ extraction is small.
- $B_d \rightarrow M^+ M^-$ decays
 - Theoretically challenging: understand the limit $m_b \rightarrow \infty$
 - Power corrections, hadronic input.
 - Tevatron measurements of $B_d \rightarrow M^+ M^-$ are competitive with B-factories.
 - I look forward to $B_s \rightarrow M^+ M^-$, ...
- Apologies. Much more B-physics at the Tevatron than in my talk...just keep listening!

Extra slides

B → D M

Decay	Br(10^{-3})	A (10^{-7} GeV)	Decay	Br(10^{-3})	A (10^{-7} GeV)
$\bar{B}^0 \rightarrow D^+ \pi^-$	2.76 ± 0.25	5.99 ± 0.27	$\bar{B}^0 \rightarrow D^{*+} \pi^-$	2.76 ± 0.21	6.06 ± 0.23
$B^- \rightarrow D^0 \pi^-$	4.98 ± 0.29	7.72 ± 0.22	$B^- \rightarrow D^{*0} \pi^-$	4.6 ± 0.4	7.50 ± 0.33
$\bar{B}^0 \rightarrow D^0 \pi^0$	0.25 ± 0.02	1.81 ± 0.08	$\bar{B}^0 \rightarrow D^{*0} \pi^0$	0.28 ± 0.05	1.95 ± 0.18
$\bar{B}^0 \rightarrow D^+ \rho^-$	7.7 ± 1.3	10.2 ± 0.9	$\bar{B}^0 \rightarrow D^{*+} \rho^-$	6.8 ± 0.9	9.10 ± 0.61
$B^- \rightarrow D^0 \rho^-$	13.4 ± 1.8	12.9 ± 0.9	$B^- \rightarrow D^{*0} \rho^-$	9.8 ± 1.7	10.5 ± 0.92
$\bar{B}^0 \rightarrow D^0 \rho^0$	0.29 ± 0.11	1.97 ± 0.37	$\bar{B}^0 \rightarrow D^{*0} \rho^0$	< 0.51	< 2.78

The Experimental Program for $\sin 2\beta_{\text{eff}}$

Mode	CP	Tot. error Belle $\mathcal{L} \sim 253 \text{ fb}^{-1}$	Tot. error BABAR $\mathcal{L} \sim 195\text{-}212 \text{ fb}^{-1}$	$\langle \Delta(\text{SM}) \rangle$ [in σ]	Error estimate at 2 ab^{-1}	Systematics	Max. central value for 5σ deviation at 2 ab^{-1}	Quality [naïve theoretical cleanliness]
ϕK^0	-1	0.34	0.26	-1.9	0.10	small	0.22	☺ ☺ ☺
$\eta' K^0$	-1	0.18	0.14	-2.6	< 0.05	small	0.45	☺ ☺ (☺)
$f_0(980) K^0$	+1	0.42	0.29	-1.3	< 0.12	Q2B	0.12	☺ ☺
$K_S K_S K^0$	± 1	0.71	0.36	-1.4	< 0.16	vertex	-0.08	☺ ☺ ☺
$K^+ K^- K^0$	$\sim +1$	0.25	0.25	-1.1	< 0.08	CP	0.31	☺ (☺)
$\pi^0 K_S$	-1	0.60	0.32	-1.4	0.13	vertex	0.07	☺
ωK^0	-1	0.66	0.36	-0.6	< 0.15	small	-0.03	(☺)
$\rho^0 K^0$	-1	-	-	?	?	Q2B	?	(☺)
ηK_S	+1	-	-	?	?	vertex	?	-
Average	-	0.39 ± 0.11	0.45 ± 0.09	-3.7	< 0.034	ok	0.53	☺ ☺

3. Lifetime measurements in $b \rightarrow s\bar{s}s$ decays

A combined fit to CP asymmetries in rare hadronic $b \rightarrow s\bar{q}q$ decays measured at BaBar and BELLE indicates a deviation from the Standard Model by 3.8σ (O. Tajima, Aspen 2005).

Consider a new CP phase σ in the $b \rightarrow s\bar{s}s$ decay. Let now $\bar{B}_s \rightarrow f_{CP+}$ denote a $b \rightarrow s\bar{s}s$ decay into a CP-even final state, e.g. $f_{CP+} = (\phi\phi)_{L=0}$. With

$$\langle f_{CP+} | B_s \rangle \propto e^{i\sigma} \quad \text{and} \quad \langle f_{CP+} | \bar{B}_s \rangle \propto -e^{-i\sigma}$$

the coefficients in

$$\Gamma[f, t] \propto |\langle f | B_L \rangle|^2 e^{-\Gamma_L t} + |\langle f | B_H \rangle|^2 e^{-\Gamma_H t}$$

read:

$$|\langle f_{CP+} | B_L \rangle|^2 \propto \frac{1 + \cos(\phi + 2\sigma)}{2}, \quad |\langle f_{CP+} | B_H \rangle|^2 \propto \frac{1 - \cos(\phi + 2\sigma)}{2}$$

$$\Gamma[f_{CP^+}, t] \propto \frac{1 + \cos(\phi + 2\sigma)}{2} e^{-\Gamma_L t} + \frac{1 - \cos(\phi + 2\sigma)}{2} e^{-\Gamma_H t}$$

For the Standard Model case $\phi = \sigma = 0$ only B_L can decay into f_{CP^+} and the lifetime measured in e.g. $(\bar{B}_s) \rightarrow (\phi\phi)_{L=0}$ determines Γ_L .

If the lifetime measured in $(\bar{B}_s) \rightarrow (\phi\phi)_{L=0}$ is longer than the one measured in $(\bar{B}_s) \rightarrow (J/\psi\phi)_{L=0}$, new physics in the $b \rightarrow s\bar{s}s$ decay amplitude is established through $\sigma \neq 0$, with the possibility of $\phi = 0$ or $\phi \neq 0$.

If the lifetime measured in $(\bar{B}_s) \rightarrow (\phi\phi)_{L=0}$ is shorter than the one measured in $(\bar{B}_s) \rightarrow (J/\psi\phi)_{L=0}$, new physics in both the $b \rightarrow s\bar{s}s$ decay amplitude and $B_s - \bar{B}_s$ mixing is established through $\sigma \neq 0$ and $\phi \neq 0$.

The same argument applies to the $b \rightarrow s\bar{u}u$ amplitude triggering $(\bar{B}_s) \rightarrow K^+K^-$ except that here a small tree amplitude is present, so that $\sigma \neq 0$ already in the Standard Model.